

# Progress in Reducing AC Losses of Bi2223 Tapes with Interfilamentary Resistive Barriers

Ryoji Inada, Yasuhiro Okumura, Akio Oota, Chengshan Li, and Pingxiang Zhang

**Abstract**—This paper presents our recent progress for the development of low-AC loss Bi2223 tapes with interfilamentary oxide barriers. For the compatibility with Bi2223 phase formation during sintering, SrZrO<sub>3</sub> was selected as barrier materials. Moreover, small amount of Bi2212 was mixed with SrZrO<sub>3</sub> to improve its ductility for cold working. Although some breakages of barrier layers still existed, the effective transverse resistivity was approximately 10 times higher than a tape with pure Ag matrix. By controlling barrier thickness, reducing a tape width below 3 mm and twisting the filaments with its length below 5 mm, coupling frequency  $f_c$  attained to 260 Hz in an AC perpendicular transverse field. Critical current densities  $J_c$  of our twisted barrier tapes were ranged in 12–15 kA/cm<sup>2</sup> at 77 K and self-field, which was 25–30 % lower than non-twisted one (= 18 kA/cm<sup>2</sup>). In our knowledge, this is the first report to achieve both  $J_c > 12$  kA/cm<sup>2</sup> and  $f_c > 250$  Hz simultaneously in a single Bi2223 tapes. Our barrier tapes showed 60-70% lower perpendicular field losses than a conventional 4 mm-width tape with fully coupled filaments at 50 mT and 50 Hz. These results are promising for remarkable improvement in AC performance for Bi2223 tapes in future.

**Index Terms**—Bi2223 tapes, interfilamentary barriers, AC loss, filament twisting, coupling frequency

## I. INTRODUCTION

At present, the Ag-sheathed (Bi,Pb)<sub>2</sub>Sr<sub>2</sub>Ca<sub>2</sub>Cu<sub>3</sub>O<sub>x</sub> (Bi2223) tapes with high critical current density ( $J_c$ ) of 40-50 kA/cm<sup>2</sup> and long length above 1 km are commercially available [1]. However, their AC losses are still too large for realization of AC power devices such as transformers, motors and cables. The large loss generation in the tape subjected to an AC transverse magnetic field in perpendicular to long axis of the tape is attributed to the electromagnetic coupling between the filaments via the Ag matrix, which has a low electric resistivity (=  $0.27 \times 10^{-8}$  Ωm) at 77 K. Particularly, due to the large aspect ratio of Bi2223 tape, both the hysteresis loss ( $Q_h$ ) in superconductor and the coupling loss ( $Q_c$ ) in matrix in a perpendicular transverse field (in perpendicular to both long axis and broader face of the tape) becomes much larger than in a parallel one (in perpendicular to long axis but parallel to

broader face of the tape), and the conditions for filament decoupling become more restrictive [2–4]. In order to reduce the filament coupling in an AC perpendicular transverse field, it is necessary – in addition to twisting the filaments with a suitable pitch length – to increase the matrix resistivity by introducing oxide layers among the filaments as resistive barriers [5–11].

Coupling frequency  $f_c$ , which is related with coupling time constant  $\tau_c$  as the expression of  $f_c = 1/2\pi\tau_c$ , is one of the important parameters to determine the AC operating condition for filament decoupling. To achieve a significant loss reduction by decoupling the filaments,  $f_c$  should at least be higher than operating frequency  $f_{op}$ . In the previous works for the development of Bi2223 tapes with oxide barriers,  $f_c$  under a perpendicular transverse field was increased above 100 Hz by introducing BaZrO<sub>3</sub> or SrZrO<sub>3</sub> barriers combined with filament twisting [6, 7]. The higher  $f_c$  of 400–500 Hz was also achieved in twisted tapes with SrZrO<sub>3</sub> with small amount of SrCO<sub>3</sub> [8–10] or Bi2212 [11] as interfilamentary barriers. However,  $J_c$  of those barrier tapes with  $f_c > 100$  Hz were limited to only 4-6 kA/cm<sup>2</sup> at 77 K and self-field.

The simultaneous achievement of  $J_c$  well above  $10^4$  A/cm<sup>2</sup> and  $f_c$  higher than several 100 Hz has crucial importance for widening the applicability of Bi2223 tapes for AC power devices. Generally, thicker barrier is preferable to maintain its continuity and increase matrix resistivity in a fully reacted tape [5]. However, it may cause not only the reduction of the oxygen diffusion paths for Bi2223 filaments but also the serious degradation of flatness for Bi2223 filaments embedded in a matrix due to the degradation for composite workability [11]. They should lead to serious  $J_c$  drops as mentioned above. Therefore, precise control in barrier thickness, tape geometries and deforming parameters during twisting and rolling process should be indispensable to obtain both high  $J_c$  and  $f_c$ .

In this paper, we report simultaneous achievement of  $J_c > 12$  kA/cm<sup>2</sup> in self-field and  $f_c > 250$  Hz in perpendicular field at 77 K for a Bi2223 tape, by controlling coating thickness of barriers before stacking and geometrical parameters such as tape widths and twist pitch lengths. Comparing the data for tapes with fully coupled filaments, loss reduction at power-grid frequency in a perpendicular transverse field was investigated at 77 K.

## II. EXPERIMENTAL

Bi2223 tapes with oxide barriers among the twisted filaments were prepared by a conventional powder-in-tube (PIT) method. SrZrO<sub>3</sub> with a mean grain size below 1 μm was used as barrier materials for its compatibility with Bi2223

Manuscript received 3 August 2010. This work was supported in part by Grant-in-Aids for Scientific Research from MEXT (No.20686020) and JSPS (No.22560270) of Japan. It was also supported in part from TEPCO Research Foundation, Research Foundation for the Electrotechnology of Chubu (No.R-20302) and DAIKO Foundation.

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superconductor. Moreover, an additional Bi2212 powder corresponding to 20wt% was mixed with SrZrO<sub>3</sub> to improve its ductility for cold working [11]. The precursor powders with a composition of Bi<sub>1.76</sub>Pb<sub>0.34</sub>Sr<sub>1.93</sub>Ca<sub>2.02</sub>Cu<sub>3.1</sub>O<sub>x</sub> were packed into a pure Ag tube with an outer diameter of 9.6 mm and a wall thickness of 0.8 mm. Then, the composite was deformed into a hexagonal cross-sectional shape by drawing, with its diagonal length of 1.8 mm. The outside surface of the monocoil wire was coated by SrZrO<sub>3</sub> + Bi2212 pastes. To obtain the sufficient workability of the composite, coating thickness of the pastes was reduced to 50–60 μm, which is approximately one half the value in our previous work [11]. After a heat treatment at 550°C in air to eliminate the organic binder in the pastes, 19-pieces of coated monocoil wire were stacked and packed into an Ag-Mg alloy tube with an outer diameter of 15.6 mm and wall thickness of 0.8 mm. The composites were drawn to the diameter of 1.33 mm and then twisted very carefully with intermediate heat treatments at 400°C in vacuum. Finally, the twisted round wires were formed into tape shapes by flat rolling, and sintered at 830–840°C with an intermediate rolling. In fully reacted tapes, cross sectional sizes are 2.7 mm × 0.23 mm and volume fractions of filaments are 23%, respectively. Twist pitch lengths  $L_t$  were measured after removing the sheath parts from final tapes by etching and ranged from 4 to 7 mm. For comparison, non-twisted tape without barriers was also fabricated by the same fabrication process.

The critical current  $I_c$  was measured in all tapes with DC four-probe method at 77 K in a self-field, with an electric field criterion of 1 μV/cm. The critical current density  $J_c$  was determined from  $I_c$  and transverse cross-sectional area of the filaments. The AC losses  $Q_m$  at 77 K in a perpendicular transverse field were measured by a saddle shaped pick-up coil and a conventional lock-in technique [12]. For the loss measurements, the lengths of tapes were fixed to 80 mm.

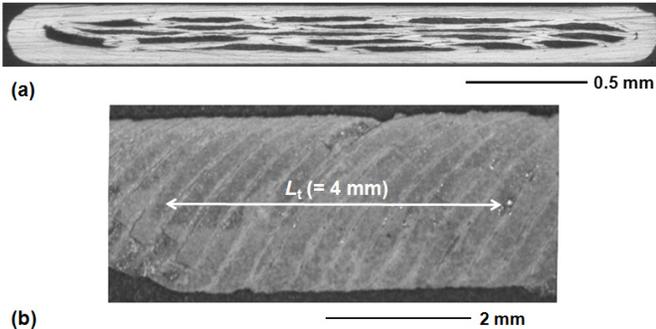


Fig. 1. (a) Transverse cross section and (b) plan view (after removing sheath part by etching) for twisted barrier tape. The size of tape section and twist pitch length are 2.7 mm × 0.24 mm and 4 mm, respectively.

### III. RESULTS AND DISCUSSION

Fig. 1 shows the transverse cross sectional view and the plan view (after removing sheath part) of twisted barrier tape with  $L_t = 4$  mm. In our previous study [11], it was confirmed that in barrier tapes with  $L_t < 10$  mm, the filaments positioned at an inner part of a tape section were distorted irregularly and physically connected each other through the broken parts of barrier layers. On the other hand, the filament shape in newly prepared barrier tape with tightly twisted filaments ( $L_t = 4$  mm)

seems to be flat. Each filament was partitioned by barrier layers and interfilamentary connections were not observed. Such fine structure would be attributed to improvements for workability of a composite by reducing coating thickness of barrier oxide and applying careful twisting. Transport critical current density  $J_c$  at 77 K and self-field for twisted tape with SrZrO<sub>3</sub> + Bi2212 barrier are shown in Fig. 2, as a function of inverse of twist pitch lengths  $L_t$ . For comparison, our previous data for barrier tapes are also plotted [11]. As can be seen,  $J_c$  of both non-twisted and twisted barrier tapes with different  $L_t$  were improved remarkably compared with previous data. For twisted barrier tapes with  $L_t < 7$  mm, their  $J_c$  values were ranged in 12–15 kA/cm<sup>2</sup> at 77 K and self-field, which was 25–30% lower than non-twisted one (= 18 kA/cm<sup>2</sup>). We consider that avoidance of irregularly distorted filaments strongly contributes to  $J_c$  improvement in twisted barrier tapes.

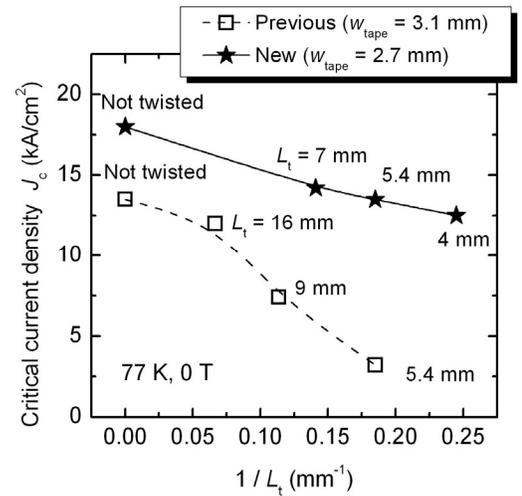


Fig. 2. Critical current densities  $J_c$  at 77 K and self-field for barrier tapes plotted against the inverse of twist pitch lengths  $L_t$ . Our previous data for barrier tapes are also shown for comparison [11].

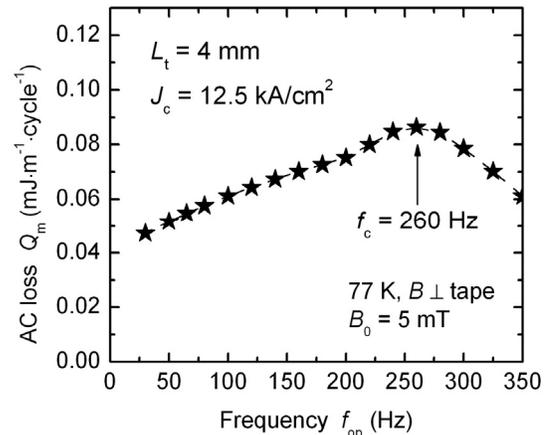


Fig. 3. Frequency dependence of AC losses  $Q_m$  per-cycle at 77 K and  $B_0 = 5$  mT for twisted barrier tape with  $L_t = 4$  mm under an AC perpendicular field.

For the tape with the shortest  $L_t = 4$  mm, AC loss properties at 77 K in a perpendicular transverse field were examined. Fig. 3 shows the frequency dependence of losses  $Q_m$  per-cycle at 77 K and fixed field amplitude  $B_0 = 5$  mT. As can be seen,  $Q_m$  data show the maximum around operating frequency  $f_{op} = 260$  Hz. This specific frequency corresponds to coupling frequency  $f_c$  at which coupling loss  $Q_c$  per-cycle included in total  $Q_m$  show the

maximum. Although the achievement for higher  $f_c$  of 400–500 Hz in barrier tapes was already reported [8–11], transport  $J_c$  of these barrier tapes with such high  $f_c$  were limited to several  $\text{kA}/\text{cm}^2$ . It should be noted that this is the first achievement for both  $J_c > 12 \text{ kA}/\text{cm}^2$  and  $f_c > 250 \text{ Hz}$  simultaneously in a single Bi2223 tape. To estimate the effective transverse resistivity  $\rho_{\perp}$  of our barrier tape, we also measured  $f_c$  for twisted tape with resistive Ag-8%Au alloy matrix and the same geometrical parameters as the barrier tape. It was found that Ag-8%Au sheath with resistivity 7–8 times higher than pure Ag at 77 K give  $f_c = 160 \text{ Hz}$ . From the extrapolation using the relation of  $f_c \propto \rho_{\perp}/L_t^2$ ,  $\rho_{\perp}$  of our barrier tape is suggested to be 10–12 times higher than a pure Ag-sheathed tape without barriers.

Perpendicular field amplitude dependence of losses at 77 K and 45 Hz are shown in Fig. 4. The data for non-twisted tapes with their tape widths ( $w_{\text{tape}}$ ) of 4 mm and 2.7 mm are also plotted as the references. These two reference tapes have no barrier layers and all filaments are electromagnetically coupled among them and behave as a single superconductor under a perpendicular field at 45 Hz. In addition, the loss values for each tape are normalized by its critical current  $I_c$  at 77 K and self-field for direct comparison among the tapes. As can be seen, the losses for twisted barrier tapes with  $L_t = 4 \text{ mm}$  are reduced by 45–60%, compared with the reference tape with  $w_{\text{tape}} = 2.7 \text{ mm}$  at  $B_0$  from 10 to 50 mT. Such remarkable loss reduction around power-grid frequency is attributed to achievement for both  $f_c > 250 \text{ Hz}$  and  $J_c > 10^4 \text{ A}/\text{cm}^2$ . In addition, the losses for the twisted barrier tapes are 60–70% lower than those for the reference tape with wider  $w_{\text{tape}} = 4 \text{ mm}$ .

To examine the loss generation mechanism of our barrier tape, Fig. 5 shows normalized loss factors  $\Gamma = \mu_0 Q_m / 2B_0^2 S_{\text{tape}}$  ( $S_{\text{tape}}$ : cross sectional area of the tape) for the barrier tape with  $L_t = 4 \text{ mm}$  at different fixed frequencies  $f_{\text{op}}$  from 30 to 105 Hz, as a function of applied field amplitude  $B_0$ . For comparison, the  $\Gamma$  curve derived from the analytical prediction of hysteresis loss  $Q_h$  for fully coupled filaments using an elliptical model approximation is also plotted [13]. Here,  $B_0$  at which  $\Gamma$  shows a maximum is defined as a parameter  $B_{\text{max}}$ .  $B_{\text{max}}$  nearly corresponds to the full penetration field  $B_p$  for the tape and changes depending on the conditions for filament coupling under AC fields. As can be seen,  $B_{\text{max}}$  for our barrier tape at different  $f_{\text{op}}$  are nearly the constant ( $\sim 1.5 \text{ mT}$ ) and lower than the calculation assuming the filaments are coupled ( $\sim 7 \text{ mT}$ ). This indicates that all filaments in the tape are decoupled and hysteresis loss component  $Q_h$  is reduced by the level for decoupled filaments below 105 Hz. At  $B_0 > 1.5 \text{ mT}$ , there are greater losses for the higher  $f_{\text{op}}$  while calculated  $\Gamma$  for coupled filaments is higher than measured ones at different  $f_{\text{op}}$ . We preliminarily confirmed that eddy current losses  $Q_e$  for a pure Ag tape with the same size as our barrier tape are much smaller than  $Q_m$  for our barrier tape at  $f_{\text{op}} < 100 \text{ Hz}$ . Therefore, it is considered that the contribution of  $Q_e$  in sheath of barrier tape is negligible and the increase of  $\Gamma$  with frequency in Fig. 5 is mainly caused by the contribution of coupling loss  $Q_c$ .

By using an effective medium approximations for the composite [14], frequency dependence of  $Q_c$  per-cycle at fixed  $B_0$  are approximately described as

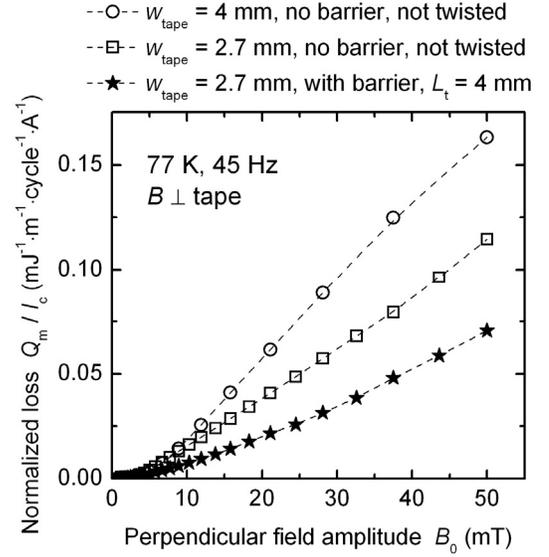


Fig. 4. Perpendicular field amplitude dependence of normalized AC losses  $Q_m/I_c$  at 77 K and 45 Hz for twisted barrier tape with  $L_t = 4 \text{ mm}$ . The data for non-twisted tapes with tape width of 4 mm and 2.7 mm are also plotted as the references. These two reference tapes don't have barrier layers and the filaments in them are fully coupled at 45 Hz.

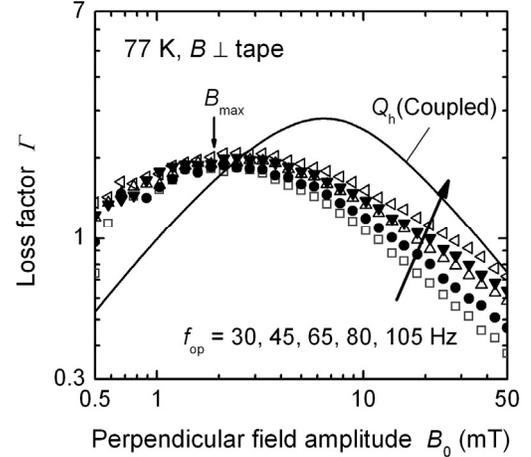


Fig. 5. Loss factor  $\Gamma = \mu_0 Q_m / 2B_0^2 S_{\text{tape}}$  for twisted barrier tape with  $L_t = 4 \text{ mm}$  at 77 K and various fixed frequency  $f_{\text{op}}$ , plotted against perpendicular field amplitude  $B_0$ . The analytical prediction of  $\Gamma$  for hysteresis loss  $Q_h$  for fully coupled filaments is also shown.

$$Q_c \propto \frac{\omega \tau_c}{1 + (\omega \tau_c)^2} \quad (1)$$

Here,  $\omega$  is equal to  $2\pi f_{\text{op}}$  and  $\tau_c = 1/2\pi f_c$  is the coupling time constant. For more precise evaluation, one should expect a skin-effect type behavior at higher frequencies, i.e.,  $Q_c$  per cycle is inversely proportional to the square root of the frequency  $f_{\text{op}} > f_c$ , which is expressed as follows [15, 16]:

$$Q_c \propto \frac{1}{\pi \sqrt{\omega \tau_c / 2}} \frac{\sinh(\pi \sqrt{\omega \tau_c / 2}) - \sin(\pi \sqrt{\omega \tau_c / 2})}{\cosh(\pi \sqrt{\omega \tau_c / 2}) + \cos(\pi \sqrt{\omega \tau_c / 2})} \quad (2)$$

Here, we used the latter expression to describe the frequency dependence of  $Q_c$  for our barrier tape. On the other hand, as mentioned above,  $Q_h$  for our barrier tape at fixed  $B_0$  is nearly independent of  $f_{\text{op}} < 110 \text{ Hz}$ . Therefore, total loss  $Q_m$  for our barrier tape is expressed as the sum of  $Q_h$  and  $Q_c$ :

$$Q_m = Q_h + \frac{q_c}{\pi \sqrt{f_{\text{op}} / 2f_c}} \frac{\sinh(\pi \sqrt{f_{\text{op}} / 2f_c}) - \sin(\pi \sqrt{f_{\text{op}} / 2f_c})}{\cosh(\pi \sqrt{f_{\text{op}} / 2f_c}) + \cos(\pi \sqrt{f_{\text{op}} / 2f_c})} \quad (3)$$

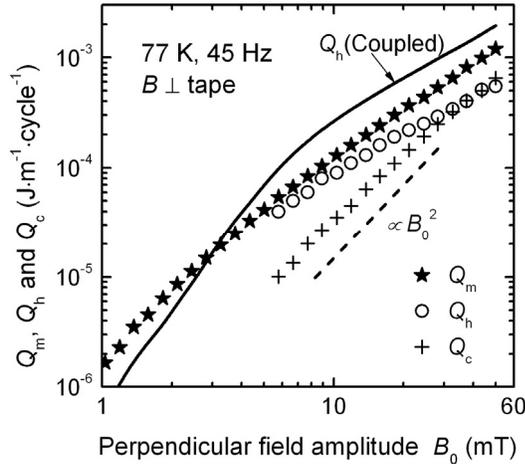


Fig. 6. Comparison of hysteresis loss  $Q_h$ , coupling loss  $Q_c$  and total loss  $Q_m$  at 77 K and 45 Hz for twisted barrier tape with  $L_t = 4$  mm, as a function of perpendicular field amplitude  $B_0$ . An analytical prediction of  $Q_h$  of an elliptical model for fully coupled filaments is also shown.

where  $q_c$  is the constant related to the magnitude of  $Q_c$  at fixed  $B_0$ . From the data shown in Fig. 3,  $f_c$  of our barrier tape was confirmed to be 260 Hz, so that  $\tau_c$  of the tape was estimated to be 0.6 ms. By fitting Eq. (3) into measured frequency dependence of total  $Q_m$  below 110 Hz, both  $Q_h$  and  $q_c$  in barrier tape can be roughly estimated. Using  $q_c$  obtained from fitting and  $f_{op}$ ,  $Q_c$  per-cycle is calculated by the second term in Eq. (3).

Fig. 6 shows the comparison of  $Q_h$  and  $Q_c$  at 45 Hz for twisted barrier tape, obtained by fitting by Eq. (3) into measurement. For comparison, total  $Q_m$  measured at 45 Hz and analytical  $Q_h$  curve for fully coupled filaments are shown [13]. As can be seen,  $Q_h$  shows linear dependence on  $B_0$  and approximately 70–80% lower than fully coupled level at  $B_0 > 10$  mT. On the other hand,  $Q_c$  is nearly proportionate to  $B_0^2$ , which is expected from analytical prediction [14–16]. At  $B_0 > 30$  mT, it is evident that the magnitude of  $Q_c$  at 45 Hz is nearly the same as  $Q_h$ . Consequently, total  $Q_m$  at  $B_0 > 30$  mT are two times higher than  $Q_h$  for decoupled filaments. This suggests that around power-grid frequency, the loss reduction of our twisted barrier tape is still limited by significant  $Q_c$  contribution. To achieve more remarkable loss reduction and also to maintain the effect in higher  $B_0$  range ( $\sim 0.1$  T), both increasing coupling frequency  $f_c$  to reduce the absolute  $Q_c$  values and improving  $J_c$  to reduce the fraction of  $Q_c$  in total  $Q_m$  should be necessary. The optimization of fabrication process and tape structure (tape width, barrier thickness and twist pitch length) is currently being studied, to improve the electromagnetic performance (both  $f_c$  and  $J_c$ ) for twisted barrier tapes.

#### IV. CONCLUSION

Our recent activities for the development of low-AC loss Bi2223 tapes with interfilamentary  $\text{SrZrO}_3 + \text{Bi2212}$  as resistive barriers were presented. By controlling coating thickness of barriers before stacking, reducing a tape width ( $< 3$  mm) and careful filament twisting, distorted filaments and their physical connections were greatly reduced in fully reacted barrier tapes with tightly twisted filaments. For the tape with twist pitch length  $L_t = 4$  mm, coupling frequency  $f_c$  exceeded

250 Hz in an AC perpendicular field. Critical current densities  $J_c$  of tightly twisted barrier tapes with twist pitches  $L_t = 4\text{--}7$  mm were ranged in 12–15 kA/cm<sup>2</sup> at 77 K and self-field. In our knowledge, this is the first achievement for both  $J_c > 12$  kA/cm<sup>2</sup> and  $f_c > 250$  Hz simultaneously in a single Bi2223 tape. The barrier tape with  $L_t = 4$  mm also showed 60–70% lower perpendicular field losses than a conventional 4 mm-width tape with fully coupled filaments at 10–50 mT and 50 Hz. These achievements are promising for remarkable improvement for AC performance of Bi2223 tapes in near future.

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